

# Variational Integrators

## for Simulation and Control

### Motivation

“Approximate the action instead of the equations of motion” A. J. Lew

**Variational Integrators** (VIs) offer many advantages

- preserve momenta, if the system has symmetries,
  - backward error analysis reveals discrete time paths as exact solutions of a nearby Hamiltonian,
  - energy error is bounded, thus excellent longtime behavior.
- Moreover, VIs are eminently suited for optimal control, since they
- result in one-step maps that can be directly utilized as discrete state-space models,
  - yield an unified computational framework for simulation of the mechanical system and optimization of the control.

### Higher Order VIs

Construction in two steps

1. approximation of the state variables in time

$$\mathbf{q}(t) \approx \mathbf{q}^d(t) = \sum_{n=0}^p M_n(t) \mathbf{q}_{k+n/p}$$

2. time-step-wise quadrature of the action-integral..

$$\begin{aligned} \Delta S &= \int_{t_k}^{t_{k+1}} L(\mathbf{q}(t), \dot{\mathbf{q}}(t), t) dt \\ &\approx \int_{t_k}^{t_{k+1}} L(\mathbf{q}^d(t), \dot{\mathbf{q}}^d(t), t) dt \\ &\approx \sum_{m=1}^g w_m L(\mathbf{q}^d(t_m), \dot{\mathbf{q}}^d(t_m), t_m) = L_d \end{aligned}$$

..and the virtual work of the nonconservative forces

$$\begin{aligned} \delta W^{\text{nc}} &= \int_{t_k}^{t_{k+1}} \mathbf{F} \cdot \delta \mathbf{q} dt \approx \int_{t_k}^{t_{k+1}} \mathbf{F} \cdot \delta \mathbf{q}^d dt \\ &\approx \sum_{m=1}^g w_m \mathbf{F}(t_m) \cdot \delta \mathbf{q}^d(t_m) = \sum_{n=0}^p \mathbf{F}_{k+n/p}^d \delta \mathbf{q}_{k+n/p}^d \end{aligned}$$

The stationarity condition of the action functional yields Discrete Euler-Lagrange-Equations (position-momentum form)

$$\begin{aligned} \mathbf{p}_k &= -D_1 L_d(\mathbf{q}_k, \mathbf{q}_{k+1/p} \dots, \mathbf{q}_{k+1}) - \mathbf{F}_k^d \\ \mathbf{0} &= D_2 L_d(\mathbf{q}_k, \mathbf{q}_{k+1/p} \dots, \mathbf{q}_{k+1}) + \mathbf{F}_{k+1/p}^d \\ &\dots \\ \mathbf{0} &= D_p L_d(\mathbf{q}_k, \mathbf{q}_{k+1/p} \dots, \mathbf{q}_{k+1}) + \mathbf{F}_{k+\frac{p-1}{p}}^d \\ \mathbf{p}_{k+1} &= D_{p+1} L_d(\mathbf{q}_k, \mathbf{q}_{k+1/p} \dots, \mathbf{q}_{k+1}) + \mathbf{F}_{k+1}^d \end{aligned}$$

where  $D_1 L_d = \frac{\partial L_d}{\partial \mathbf{q}_k}$ ,  $D_2 L_d = \frac{\partial L_d}{\partial \mathbf{q}_{k+1/p}}$ ,  $\dots$ ,  $D_{p+1} L_d = \frac{\partial L_d}{\partial \mathbf{q}_{k+1}}$ .

**References:** J. Marsden, T. Murphey, S. Leyendecker

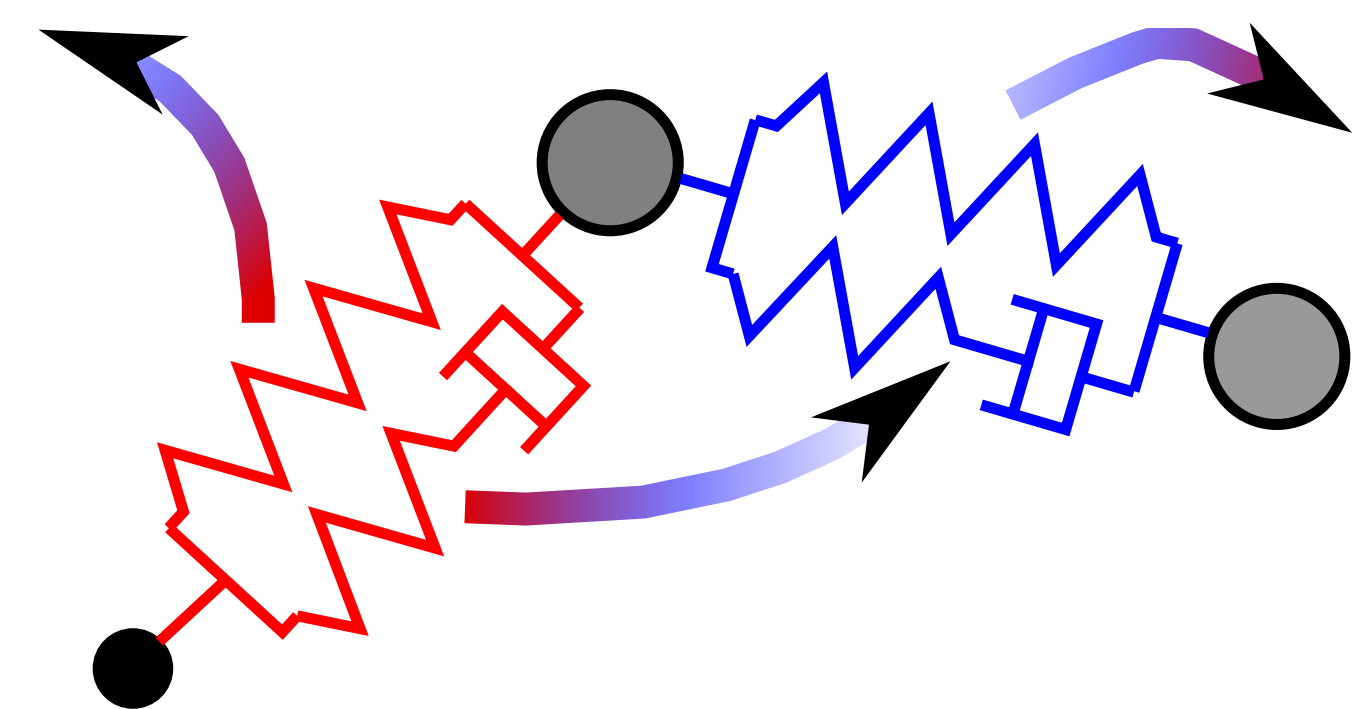
### Thermo-viscoelastic Systems

Variational formulation

$$\delta \int_{t_0}^{t_1} (T^* - \psi) dt + \int_{t_0}^{t_1} \delta W^{\text{nc}} dt = 0$$

where may be inserted the exemplaric expressions

kinetic coenergy	$T^* = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2)$
elastic strain energy	$\psi_e = \frac{K}{2l_0^2} (l - l_0)^2$
thermoelastic coupling	$\psi_{te} = -\beta (\vartheta - \vartheta_r) \frac{l - l_0}{l_0}$
heat capacity	$\psi_t = -\frac{k}{2\vartheta_r} (\vartheta - \vartheta_r)^2$
heat flux/source	$\delta W_t^{\text{nc}} = \dot{s} \delta \alpha$
internal dissipation	$\delta W_v^{\text{nc}} = F_v \delta v$



thermo-viscoelastic double pendulum as example

**References:** I. Romero, G. Maugin, A. Bertram

### Outlook

Future work is related with

**constraints**

$$\delta S_c = \delta \int_{t_0}^{t_1} L(q, \dot{q}, t) - \sum_{i=1}^{N_c} \lambda_i \varphi_i dt + \int_{t_0}^{t_1} \delta W^{\text{nc}} dt = 0$$

**electromechanical systems**

$$L = T^* + W_m^* - V - W_e$$

**optimal control**

$$\text{cost} \quad J = \int_{t_0}^{t_1} \ell(t, z, u) dt \quad \text{dynamics} \quad \dot{z} = f(z, u)$$

**References:** T. Murphey, D. Liberzon, S. Ober-Blöbaum

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  - Variational Integrators
  - Dynamics and Control